# On periodic traveling waves of the Camassa-Holm equation 

## Anna Geyer

joint work with Jordi Villadelprat


Theoretical and Computational Aspects of Nonlinear Surface Waves

BIRS
Oct 30 - Nov 42016

## The Camassa-Holm equation

$$
\begin{equation*}
u_{t}-u_{t x x}+2 k u_{x}+3 u u_{x}=2 u_{x} u_{x x}+u u_{x x x} \tag{CH}
\end{equation*}
$$

- arises as a shallow water approximation of the Euler equations for inviscid incompressible homogeneous fluids.
- $u=u(t, x)$ represents the water's free surface and $k \in \mathbb{R}$ is a parameter related to the critical shallow water speed.
- The CH equation is completely integrable and
- models wave breaking.
R. Camassa and D. Holm, An integrable shallow water equation with peaked solitons, Phys. Rev. Lett. 71 (1993)


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Traveling Wave Solutions: $u(x, t)=\varphi(x-c t)$

$$
\begin{equation*}
\varphi^{\prime \prime}(\varphi-c)+\frac{\left(\varphi^{\prime}\right)^{2}}{2}+r+(c-2 k) \varphi-\frac{3}{2} \varphi^{2}=0 \tag{1}
\end{equation*}
$$

where $c$ is the wave speed and $r \in \mathbb{R}$ is an integration constant.
J. Lenells, Traveling wave solutions of the Camassa-Holm equation, JDE '05.

## Camassa-Holm equation

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\end{equation*}
$$

where $c$ is the wave speed and $r \in \mathbb{R}$ is an integration constant. We will concentrate on smooth periodic TWS.

$\lambda \ldots$ wave length a ... wave height

## Proposition (Waves $\longleftrightarrow$ Orbits)

- $\varphi$ is a smooth periodic solution of (2)

$$
\varphi^{\prime \prime}(\varphi-c)+\frac{\left(\varphi^{\prime}\right)^{2}}{2}+r+(c-2 k) \varphi-\frac{3}{2} \varphi^{2}=0
$$

if and only if $(w, v)=\left(\varphi-c, \varphi^{\prime}\right)$ is a periodic orbit of the planar system

$$
\left\{\begin{array}{l}
w^{\prime}=v  \tag{3}\\
v^{\prime}=-\frac{A^{\prime}(w)+\frac{1}{2} v^{2}}{w}
\end{array}\right.
$$

where

$$
A(w):=\alpha w+\beta w^{2}-\frac{1}{2} w^{3}
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with $\alpha:=r-2 k c-\frac{1}{2} c^{2}$ and $\beta:=-(c+k)$.

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- Every periodic orbit of (3) belongs to the period annulus $\mathscr{P}$ of a center, which exists if and only if $-2 \beta^{2}<3 \alpha<0$.


## Observations:

periodic solution $\varphi$ wave lenght $\lambda$ of $\varphi$ wave height $\boldsymbol{a}$ of $\varphi=$
$\left\{\varphi_{a}\right\}_{a \in\left(0, a_{M}\right)}$
$\longleftrightarrow \quad$ periodic orbit $\gamma$
$=\quad$ period $\boldsymbol{T}$ of $\gamma$
$=\quad \ell(h)$, where $\ell$ is an analytic diffeo with $\ell\left(h_{0}\right)=0$.
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$\left\{\varphi_{a}\right\}_{a \in\left(0, a_{M}\right)} \longleftrightarrow \quad\left\{\gamma_{h}\right\}_{h \in\left(h_{0}, h_{1}\right)}$
Consequence: $\lambda(a)$ is a well-defined function

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\begin{aligned}
& \lambda:\left(0, a_{M}\right) \longrightarrow \mathbb{R}^{+} \\
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Deduce qualitative properties of the function $\lambda$ from those of the period function $\mathbf{T}$.

Result: $\lambda(a)$ is either unimodal or monotonous.

Theorem (A.G. \& J. Villadelprat)
Given $c, k, c \neq-k$, there exist real numbers $r_{1}<r_{b_{1}}<r_{b_{2}}<r_{2}$ such that the Camassa-Holm equation

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has smooth periodic TWS $\varphi(x-c t)$ satisfying

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- If $r \in\left(r_{1}, r_{b_{1}}\right]$, then $\lambda(a)$ is monotonous increasing.
- If $r \in\left(r_{b_{1}}, r_{b_{2}}\right)$, then $\lambda(a)$ has a unique critical point (maximum).
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The bifurcation values are $r_{1}=-\frac{2}{3}\left(k-\frac{1}{2} c\right)^{2}, r_{2}=c\left(\frac{1}{2} c+2 k\right)$, and
$r_{b_{1}}=k\left(c-\frac{1}{2} k\right), r_{b_{2}}=\frac{\sqrt{6}-3}{12}(3 k \sqrt{6}+2 c+8 k)(k \sqrt{6}-2 c-2 k)$

- For the Degasperis-Procesi equation

$$
u_{t}+2 k u_{x}+4 u u_{x}-u_{t x x}=3 u_{x} u_{x x}+u u_{x x x}
$$

the TWS satisfy an equation of the form

$$
\varphi^{\prime \prime}(\varphi-c)+\left(\varphi^{\prime}\right)^{2}+r+(c-2 k) \varphi-2 \varphi^{2}=0
$$

and the wave length is qualitatively the same as for CH .

The wave length $\lambda(a)$ is monotonous for periodic TWS of

- Korteweg-de Vries and BBM equation:

$$
u_{t}+u_{x}+\frac{3}{2} u u_{x}+\alpha u_{x x x}+\beta u_{t x x}=0
$$

with $\alpha, \beta \in \mathbb{R}$, whose TWS satisfy an equation of the form

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\varphi^{\prime \prime}+f(\varphi)=0
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where $f$ is quadratic.
$\rightsquigarrow$ use monotonicity criteria by [Chicone, '87] or [Schaaf, '85].

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- generalized reduced Ostrovsky:

$$
\left(u_{t}+u^{p} u_{x}\right)_{x}=u
$$

with $p \in \mathbb{N}$, whose TWS satisfy an equation of the form

$$
\varphi^{\prime \prime}\left(\varphi^{p}-c\right)+p \varphi^{(p-1)}\left(\varphi^{\prime}\right)^{2}-\varphi=0
$$

$\rightsquigarrow$ classical monotonicity criteria do not apply.

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H(x, y)=V(x)+\frac{1}{2} y^{2}
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[F. Mañosas, J. Villadelprat, JDE '09]

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- potential systems:

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[F. Mañosas, J. Villadelprat, JDE '09]

- systems with quadratic-like centers:

$$
H(x, y)=A(x)+B(x) y+C(x) y^{2}
$$

[A. Garijo, J. Villadelprat, JDE '14]

Consider a Hamiltonian differential system of the form

$$
\left\{\begin{array}{l}
x^{\prime}=y, \\
y^{\prime}=-V^{\prime}(x),
\end{array} \quad H(x, y)=\frac{y^{2}}{2}+V(x)\right.
$$

where $V(x)$ is a quadratic potential. For the period function

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T(h)=\int_{\gamma_{h}} \frac{d x}{y}=\sqrt{2} \int_{x_{l}}^{x_{r}} \frac{d x}{\sqrt{h-V(x)}}
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where $m_{k}$ are defined recursively using the Hamiltonian. Then,
" critical periods $=\operatorname{zeros}\left(T^{\prime}(h)\right) \leq \operatorname{zeros}\left(m_{k}\right)=: n$, if $n<k$."
$\rightsquigarrow$ Chebyshev criterion for Abelian integrals [Grau, Mañosas, Villadelprat, '11]

Consider an analytic differential system satisfying
The system has a center at the origin, an analytic first integral of the form
(H)

$$
H(x, y)=A(x)+B(x) y+C(x) y^{2} \text { with } A(0)=0,
$$

and its integrating factor $K$ depends only on $x$.

## Theorem (Garijo \& Villadelprat, 2014)

Under hypotheses $(\mathbf{H})$ let $\mu_{0}=-1$ and define for $i \geqslant 1$

$$
\mu_{k}:=\left(\frac{1}{2}+\frac{1}{2 k-3}\right) \mu_{k-1}+\frac{\sqrt{|C|} V}{(2 k-3) K}\left(\frac{K \mu_{k-1}}{\sqrt{|C|} V^{\prime}}\right)^{\prime} \text { and } \ell_{k}:=\frac{k}{\sqrt{|C|} V^{\prime}} \mu_{k}
$$

If the number of zeros of $\mathscr{B}_{\sigma}\left(\ell_{k}\right)$ on $\left(0, x_{r}\right)$, counted with multiplicities, is $n \geqslant 0$ and it holds that $k>n$, then the number of critical periods of the center at the origin, counted with multiplicities, is at most $n$.
$\mathscr{B}_{\sigma}\left(\ell_{k}\right)$ denotes the $\sigma$-odd part of $\ell_{k}$ for an involution $\sigma$ defined in terms of $H$.

The system

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\left\{\begin{array}{l}
w^{\prime}=v, \\
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has the first integral

$$
H(w, v):=A(w)+\frac{w}{2} v^{2}
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To apply the criterion of [GaVi14], our system has to satisfy the following hypotheses:

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## Period Annuli

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\begin{cases}x^{\prime}=y, & A(x)=\frac{1}{2} x^{2}-x^{3}  \tag{4}\\ y^{\prime}=-\frac{A^{\prime}(x)+y^{2}}{2(x+\vartheta)}, & \text { where } \\ \vartheta:=\frac{1}{6}\left(\frac{2}{\sqrt{4+\frac{6 \alpha}{\beta^{2}}}}-1\right)\end{cases}
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## Proposition (1)

If $\vartheta \geqslant \frac{1}{6}$, then the period function of the center of system (4) is monotonous increasing.

- Apply the criterion in [GaVi14] to deduce monotonicity.



## Proposition (2)

For $\vartheta<\frac{1}{6}$ the period function of the center of (4) is either monotonous decreasing for $\vartheta \in\left(0, \vartheta_{1}\right]$ or unimodal
$\vartheta \in\left(\vartheta_{1}, 1 / 6\right)$, where $\vartheta_{1}=-\frac{1}{10}+\frac{1}{15} \sqrt{6}$.

- Apply criterion in [GaVi14] to obtain an upper bound for the critical periods.
- Compute the first period constants.
- Determine the sign of $T^{\prime}(h)$ for $h \approx h_{m}$.



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To apply the criterion, we study the number of roots of $\mathscr{B}_{\sigma}\left(\ell_{i}\right)$.

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| Intervals | $\left(0, \vartheta_{0}\right)$ | $\left(\vartheta_{0}, \vartheta_{1}\right)$ | $\left(\vartheta_{1}, 1 / 6\right)$ |
| :--- | :---: | :---: | :---: |
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|  |  |  |  |
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| $\Downarrow$ [GaVi14], THM A |  |  |  |
| Period function $T(h)$ | monot. | $\leq 1$ crit. per. | $\leq 2$ crit. per. |

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| $\Downarrow h \approx 0, h \approx h_{m}$ |  |  |  |
| Period function $T(h)$ | monot. decreasing |  |  |

- $T^{\prime}(h) \lessgtr 0$ for $\vartheta \lessgtr \vartheta_{1}$ near $h=0$.
- $T^{\prime}(h) \longrightarrow-\infty$ as $h \rightarrow h_{m}$.

Sketch of the graph of the period function $T(h)$ of the center of (4):


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$$
0<\vartheta \leq \vartheta_{1}
$$

$\vartheta_{1}<\vartheta<\frac{1}{6}$
$\vartheta \geq \frac{1}{6}$
Proposition (2)
Proposition (1)

$$
T_{0}=2 \pi \sqrt{2 \vartheta}, \quad T_{1}=2 \ln \left(\frac{\sqrt{(2 \vartheta+1)(1-6 \vartheta)}}{1+6 \vartheta-4 \sqrt{\vartheta(1+3 \vartheta)}}\right)>0
$$

## Summary


$0<\vartheta \leq \vartheta_{1}$

## Theorem

Given $c \neq-k$, there exist real numbers $r_{1}<r_{b_{1}}<r_{b_{2}}<r_{2}$ such that the Camassa-Holm equation has smooth periodic TWS satisfying

$$
\varphi^{\prime \prime}(\varphi-c)+\frac{\left(\varphi^{\prime}\right)^{2}}{2}+r+(c-2 k) \varphi-\frac{3}{2} \varphi^{2}=0
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if and only if $r \in\left(r_{1}, r_{2}\right)$. The set of smooth periodic
TWS form a continous family $\left\{\varphi_{a}\right\}_{a}$ parametrized by the wave height $a$.
The wave length $\lambda=\lambda(a)$ of $\varphi_{a}$ satisfies the following:

- If $r \in\left(r_{1}, r_{b_{1}}\right]$, then $\lambda(a)$ is monotonous increasing.
- If $r \in\left(r_{b_{1}}, r_{b_{2}}\right)$, then $\lambda(a)$ has a unique critical point (maximum).
- If $r \in\left[r_{b_{2}}, r_{2}\right)$, then $\lambda(a)$ is monotonous decreasing.

