

# The Conformal Method Gives a Poor Parameterization of Initial Data

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# Outline

- 1 Introduction
- 2 (Mostly Numerical) Results
- 3 Conclusions

# Einstein Constraint Equations

- Given a (vacuum, 3+1) spacetime, can take a space-like slice.  
It must satisfy

$$R + (\operatorname{tr}K)^2 - |K|^2 = 0$$

$$\operatorname{div}K - \nabla(\operatorname{tr}K) = 0.$$

- If these are satisfied, can evolve a spacetime.  
(Choquet-Bruhat '52)
- Can we parameterize all initial data?

# The Conformal Method

Main idea: split determined data from freely specifiable data.

Given  $(M^3, g, \sigma, \tau, N)$ , solve

$$8\Delta\phi = R\phi + \frac{2}{3}\tau^2\phi^5 - \left|\sigma + \frac{LW}{2N}\right|^2 \phi^{-7}$$

$$\operatorname{div} \frac{LW}{2N} = \frac{2}{3}\phi^6 d\tau$$

for  $(\phi, W)$ . Then

$$\gamma = \phi^4 g$$

$$K = \phi^{-2}(\sigma + LW/2N) + \frac{\tau}{3}\phi^4 g$$

solve the Einstein constraint equations.

$\tau$  is the mean curvature, and controls the coupling.

# Compact Solvability Results

Table: Constant mean curvature (CMC) solvability (Isenberg '95)

	$\tau = 0, \sigma \equiv 0$	$\tau = 0, \sigma \neq 0$	$\tau \neq 0, \sigma \equiv 0$	$\tau \neq 0, \sigma \neq 0$
$Y(g) > 0$	None	Unique	None	Unique
$Y(g) = 0$	"Constants"	None	None	Unique
$Y(g) < 0$	None	None	Unique	Unique

- Thus, straightforward to parameterize CMC initial data.
- Conjecture was that solvability is the same, even if not CMC.
- Small caveat: must be able to solve a prescribed scalar curvature problem if  $Y(g) < 0$ .

# Signs of Weakness

“Near-CMC” if  $d\tau$  is small compared to  $\tau > 0$ .

Table: Near-CMC Solvability, Conjectured Solvability

	$\tau \neq 0, \sigma \equiv 0$	$\tau \neq 0, \sigma \neq 0$
$Y(g) > 0$	None	Unique
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- Maxwell '11:  $T^3$  symmetric data,  $Y(g) = 0, \sigma \neq 0$  for a  $\tau$  that changes sign: Non-existence and Non-uniqueness .
- The-Cang '15:  $Y(g) > 0, \sigma$  with limited support,  $\tau > 0$ ,  $L(d\tau/\tau) \leq (d\tau/\tau)^2$ . Non-existence and Non-uniqueness.  $\sigma \equiv 0$  solutions.

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Table: Observed far-from-CMC Solvability,  $\tau > 0$

	$\tau \neq 0, \sigma \equiv 0$	$\tau \neq 0, \sigma \neq 0$
$Y(g) > 0$	Existence Non-uniqueness	Non-existence Non-uniqueness
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# Introduction to Results

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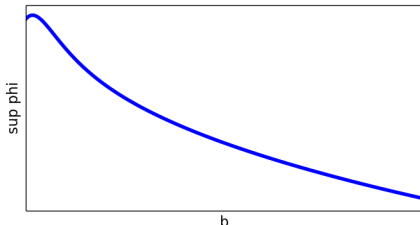
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# Methodology

- Bifurcation theory attempts to describe the family of solutions as a parameter(s) change.
- AUTO is a program for exploring ODE bifurcations numerically.
- We reduced the conformal method to an ODE by symmetry, e.g. on  $S^1 \times S^2$ .
- Usually used  $\tau = b\xi^a$ , with  $\xi > 0$ ,  $\sup \xi = 1$ .
  - $b$  gives size
  - $a$  gives "distance" from CMC, since  $d\tau^a/\tau^a = a d\tau/\tau$ .

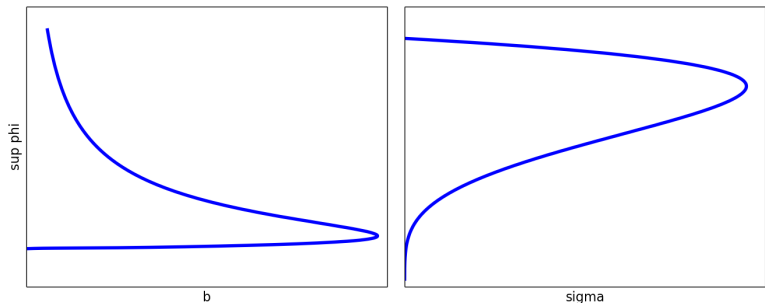


# $S^1 \times S^2$ , $S^1$ dependent ( $Y(g) > 0$ )

- $\tau = b\xi^a$ ,  $\xi > 0$ . Unique solutions for all  $b$  and  $a$ .
- Theorem: There are no  $S^1$  dependent solutions to the “limit equation” (which suggests existence/uniqueness is generic.)
- Also, no  $\tau$ 's that satisfy The-Cang's conditions.

# $S^1 \times S^2$ , latitude dependent ( $Y(g) > 0$ )

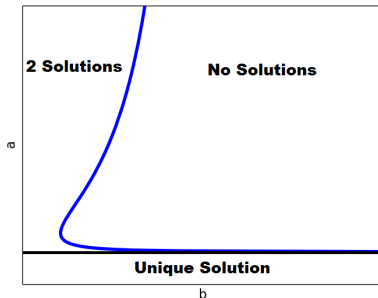
For all  $\tau = b\zeta^a$  we tested, get same generic picture for “large”  $a$ .



- Agrees with The-Cang's results.
- However, none of his conditions seem to be necessary.

# $S^1 \times S^2$ , latitude dependent ( $Y(g) > 0$ )

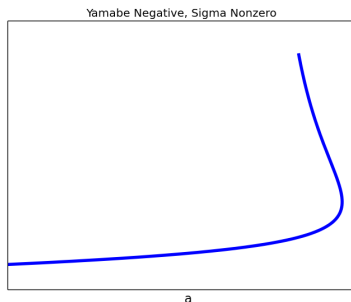
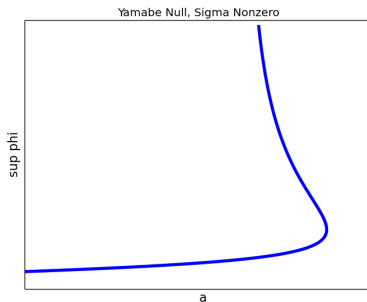
For all  $\tau = b\xi^a$  we tested, get same generic existence plot.



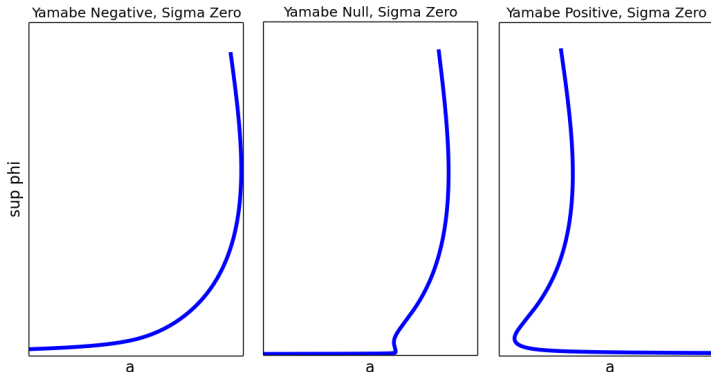
- Why does this bend? Why does larger  $a$  allow solutions?
- Unique solutions below a horizontal line, which defines “near-CMC”. Numerically related to “the” solution of the limit equation. What is this value?

## $S^2 \times H^2$ , latitude dependent

- As expected, existence and uniqueness for  $S^1 \times H^2$ ,  $\tau = b\xi^a$ .
- Instead we look at  $S^2 \times H^2$ , with a compactified  $H^2$ . Yamabe class changes as the size of  $S^2$  changes.
- Same picture for  $Y(g) > 0$  data.
- For  $\sigma \neq 0$ , similar results for  $Y(g) = 0$ ,  $Y(g) < 0$  data.



# $S^2 \times H^2$ , latitude dependent



- Can find a fold for all Yamabe classes if  $\sigma \equiv 0$ .



# The Conformal Method is Complicated

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- The conjectured solvability was very wrong for far-from-CMC data.
- If you want a simple 1-1 correspondence between specified data and initial data, the conformal method doesn't work.
- Much of this picture has not been proven analytically.

## Why do folds appear?

$$8\Delta\phi = R\phi + \frac{2}{3}\tau^2\phi^5 - \left|\sigma + \frac{LW}{2N}\right|^2\phi^{-7}$$
$$\operatorname{div}\frac{LW}{2N} = \frac{2}{3}\phi^6 d\tau$$

- $\phi^5$  is critical exponent from e.g. the Yamabe problem.
- $\tau^2$  is good sign, but the  $LW$  term is a bad sign.
- When far-from-CMC, the bad sign apparently wins out.

### Theorem (Premoselli '14)

Consider the case where  $\tau^2$  is replaced by  $\frac{2}{3}\tau_0^2 - 2\Lambda < 0$ , and  $\sigma$  by  $a\sigma$ . Then there is an  $a_* > 0$  such that

- if  $a > a_*$  there are no solutions.
- if  $a = a_*$  there is exactly one solution.
- if  $0 < a < a_*$  there are at least two solutions.

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## What else could we try?

- Examine variations of the conformal method.
- Do all (maximal) spacetimes have CMC slices? No.
- Perhaps generic spacetimes have CMC slices. Or perhaps near-CMC slices?
- Maybe there are straightforward conditions required for no CMC slices? (e.g., the initial manifold doesn't allow a Yamabe non-negative metric)

Thank you!