

# SDE limits for transfer matrices with hyperbolic channels and limiting eigenvalue processes

Christian Sadel<sup>1</sup>  
joint work with Bálint Virág<sup>2</sup>

<sup>1</sup> University of British Columbia, Vancouver

<sup>2</sup> University of Toronto

work in progress

Banff - 31 Oct 2013

# Random Schrödinger operators on strips

- Consider the random Schrödinger operator

$$(H_\lambda \psi)(n) = \psi(n+1) + \psi(n-1) + A\psi(n) + \lambda V(n)\psi(n)$$

where  $\psi = (\psi(n))_{n \in \mathbb{Z}} \in \ell^2(\mathbb{Z}) \otimes \mathbb{C}^d$ ,  $\lambda \in \mathbb{R}$  is a small coupling constant,  $A, V(n) \in \text{Her}(d)$ .

- The  $V(n)$  are i.i.d. random matrices with  $\mathbb{E}(V(n)) = \mathbf{0}$ , such that  $\mathbb{E}(\|V(n)\|^{6+\epsilon}) < \infty$ .
- If  $A$  is the adjacency matrix of a finite graph  $\mathbb{G}$  and the  $V(n)$  are diagonal matrices with i.i.d. entries along the diagonal, then this corresponds to the Anderson model on the product graph  $\mathbb{Z} \times \mathbb{G}$ .

# Random Schrödinger operators on strips

- Consider the random Schrödinger operator

$$(H_\lambda \psi)(n) = \psi(n+1) + \psi(n-1) + A\psi(n) + \lambda V(n)\psi(n)$$

where  $\psi = (\psi(n))_{n \in \mathbb{Z}} \in \ell^2(\mathbb{Z}) \otimes \mathbb{C}^d$ ,  $\lambda \in \mathbb{R}$  is a small coupling constant,  $A, V(n) \in \text{Her}(d)$ .

- The  $V(n)$  are i.i.d. random matrices with  $\mathbb{E}(V(n)) = \mathbf{0}$ , such that  $\mathbb{E}(\|V(n)\|^{6+\epsilon}) < \infty$ .
- If  $A$  is the adjacency matrix of a finite graph  $\mathbb{G}$  and the  $V(n)$  are diagonal matrices with i.i.d. entries along the diagonal, then this corresponds to the Anderson model on the product graph  $\mathbb{Z} \times \mathbb{G}$ .

- Solving  $H_\lambda \psi = E\psi$  leads to  $\begin{pmatrix} \psi_{n+1} \\ \psi_n \end{pmatrix} = \mathcal{T}_{\lambda,n}^E \begin{pmatrix} \psi_n \\ \psi_{n-1} \end{pmatrix}$  where

$$\mathcal{T}_{\lambda,n}^E = \begin{pmatrix} E\mathbf{1} - A - \lambda V(n) & -\mathbf{1} \\ \mathbf{1} & \mathbf{0} \end{pmatrix}$$

are called transfer matrices. We call  $\mathcal{T}_0^E = \mathcal{T}_{0,n}^E$  the unperturbed or free transfer matrix.

- Let  $E$  be an energy such that  $E\mathbf{1} - A$  has no eigenvalue  $\pm 2$ .
- Let  $(\varphi_j)_{j=1}^d$  be an orthonormal set of eigenvectors of  $A$ . We call  $\varphi_j$  an elliptic channel at energy  $E$  if it corresponds to an eigenvalue of  $E\mathbf{1} - A$  with absolute values  $< 2$ , and we call it a hyperbolic channel if it corresponds to an eigenvalue of  $E\mathbf{1} - A$  of absolute value  $> 2$ .

- Solving  $H_\lambda \psi = E\psi$  leads to  $\begin{pmatrix} \psi_{n+1} \\ \psi_n \end{pmatrix} = \mathcal{T}_{\lambda,n}^E \begin{pmatrix} \psi_n \\ \psi_{n-1} \end{pmatrix}$  where

$$\mathcal{T}_{\lambda,n}^E = \begin{pmatrix} E\mathbf{1} - A - \lambda V(n) & -\mathbf{1} \\ \mathbf{1} & \mathbf{0} \end{pmatrix}$$

are called transfer matrices. We call  $\mathcal{T}_0^E = \mathcal{T}_{0,n}^E$  the unperturbed or free transfer matrix.

- Let  $E$  be an energy such that  $E\mathbf{1} - A$  has no eigenvalue  $\pm 2$ .
- Let  $(\varphi_j)_{j=1}^d$  be an orthonormal set of eigenvectors of  $A$ . We call  $\varphi_j$  an elliptic channel at energy  $E$  if it corresponds to an eigenvalue of  $E\mathbf{1} - A$  with absolute values  $< 2$ , and we call it a hyperbolic channel if it corresponds to an eigenvalue of  $E\mathbf{1} - A$  of absolute value  $> 2$ .

# elliptic and hyperbolic channels

- Each elliptic channel  $\varphi_j$  gives a pair of eigenvalues  $e^{\pm ik}$ ,  $k \in (0, \pi)$  of  $\mathcal{T}_0^E$  with eigenvectors  $\begin{pmatrix} e^{\pm ik} \varphi_j \\ \varphi_j \end{pmatrix}$  corresponding to one left and one right moving wave (extended eigenstate) of  $H_0$ .
- Each hyperbolic channel gives a pair of eigenvalues  $\gamma^{\pm 1}$ ,  $|\gamma| > 1$ , for  $\mathcal{T}_0^E$ .
- If at an energy  $E$  one has  $d_e$  elliptic and  $d_h$  hyperbolic channels,  $d_e + d_h = d$ , then the multiplicity of the spectrum of  $H_0$  at  $E$  is  $2d_e$  (there are  $d_e$  overlapping bands at  $E$  for each of which one has one right and one left moving extended eigenstate)
- We want to describe the Markov process of the transfer matrix from 1 to  $n$ ,

$$\mathcal{T}_{\lambda, [1, n]}^E = \mathcal{T}_{\lambda, n}^E \mathcal{T}_{\lambda, n-1}^E \cdots \mathcal{T}_{\lambda, 2}^E \mathcal{T}_{\lambda, 1}^E$$

in a critical scaling limit  $n \sim \lambda^{-2}$ ,  $\lambda \rightarrow 0$ ,  $n \rightarrow \infty$ .

- Each elliptic channel  $\varphi_j$  gives a pair of eigenvalues  $e^{\pm ik}$ ,  $k \in (0, \pi)$  of  $\mathcal{T}_0^E$  with eigenvectors  $\begin{pmatrix} e^{\pm ik} \varphi_j \\ \varphi_j \end{pmatrix}$  corresponding to one left and one right moving wave (extended eigenstate) of  $H_0$ .
- Each hyperbolic channel gives a pair of eigenvalues  $\gamma^{\pm 1}$ ,  $|\gamma| > 1$ , for  $\mathcal{T}_0^E$ .
- If at an energy  $E$  one has  $d_e$  elliptic and  $d_h$  hyperbolic channels,  $d_e + d_h = d$ , then the multiplicity of the spectrum of  $H_0$  at  $E$  is  $2d_e$  (there are  $d_e$  overlapping bands at  $E$  for each of which one has one right and one left moving extended eigenstate)
- We want to describe the Markov process of the transfer matrix from 1 to  $n$ ,

$$\mathcal{T}_{\lambda, [1, n]}^E = \mathcal{T}_{\lambda, n}^E \mathcal{T}_{\lambda, n-1}^E \cdots \mathcal{T}_{\lambda, 2}^E \mathcal{T}_{\lambda, 1}^E$$

in a critical scaling limit  $n \sim \lambda^{-2}$ ,  $\lambda \rightarrow 0$ ,  $n \rightarrow \infty$ .

- Each elliptic channel  $\varphi_j$  gives a pair of eigenvalues  $e^{\pm ik}$ ,  $k \in (0, \pi)$  of  $\mathcal{T}_0^E$  with eigenvectors  $\begin{pmatrix} e^{\pm ik} \varphi_j \\ \varphi_j \end{pmatrix}$  corresponding to one left and one right moving wave (extended eigenstate) of  $H_0$ .
- Each hyperbolic channel gives a pair of eigenvalues  $\gamma^{\pm 1}$ ,  $|\gamma| > 1$ , for  $\mathcal{T}_0^E$ .
- If at an energy  $E$  one has  $d_e$  elliptic and  $d_h$  hyperbolic channels,  $d_e + d_h = d$ , then the multiplicity of the spectrum of  $H_0$  at  $E$  is  $2d_e$  (there are  $d_e$  overlapping bands at  $E$  for each of which one has one right and one left moving extended eigenstate)
- We want to describe the Markov process of the transfer matrix from 1 to  $n$ ,

$$\mathcal{T}_{\lambda, [1, n]}^E = \mathcal{T}_{\lambda, n}^E \mathcal{T}_{\lambda, n-1}^E \cdots \mathcal{T}_{\lambda, 2}^E \mathcal{T}_{\lambda, 1}^E$$

in a critical scaling limit  $n \sim \lambda^{-2}$ ,  $\lambda \rightarrow 0$ ,  $n \rightarrow \infty$ .



# SDE limit for $\mathcal{T}_0^E$ having only elliptic channels

Recall:

$$\mathcal{T}_{\lambda,[1,n]}^E = \mathcal{T}_{\lambda,n}^E \mathcal{T}_{\lambda,n-1}^E \cdots \mathcal{T}_{\lambda,2}^E \mathcal{T}_{\lambda,1}^E$$

## Theorem (Valko, Virag; Bachmann, de Roeck)

Let  $A$  and  $E$  be such that there are only elliptic channels (i.e.  $\mathcal{T}_0^E$  is conjugated to a unitary matrix), and consider the process  $X_{\lambda,n} = (\mathcal{T}_0^E)^{-n} \mathcal{T}_{\lambda,[1,n]}^E$ . Then, in distribution for  $n \rightarrow \infty$

$$X_{\frac{1}{\sqrt{n}},[tn]} \implies X_t$$

where  $X_t$  satisfies a SDE (stochastic differential equation) of the form

$$dX_t = d\mathcal{B}_t X_t$$

where  $d\mathcal{B}_t$  is a matrix-Brownian motion with certain variances and covariances of its entries.

## Applications in the pure one-dimensional model:

- Kritchevski, Valko and Virag studied the eigenvalue statistics in this critical scaling.
- Rifkind and Virag studied distribution of shape of eigenfunctions in this scaling limit

## Application on strip models:

- Valko and Virag obtained GOE limiting statistics for certain sequences of modified Anderson models on long boxes
- Bachmann and De Roeck discussed relations from Random Matrix Theory to the Anderson model and the DMPK equation
- In both papers the Anderson model on a strip is slightly modified by scaling down the vertical Laplacian to ensure that one has an energy interval around 0 such that for all these energies there are only has elliptic channels.
- **If you want to treat an energy interval for the honest Anderson model in a limit to infinite width or if you want to treat all energies in the spectrum of  $H_0$  in this critical limit, then one has to deal with hyperbolic channels.**

## Applications in the pure one-dimensional model:

- Kritchevski, Valko and Virag studied the eigenvalue statistics in this critical scaling.
- Rifkind and Virag studied distribution of shape of eigenfunctions in this scaling limit

## Application on strip models:

- Valko and Virag obtained GOE limiting statistics for certain sequences of modified Anderson models on long boxes
- Bachmann and De Roeck discussed relations from Random Matrix Theory to the Anderson model and the DMPK equation
- In both papers the Anderson model on a strip is slightly modified by scaling down the vertical Laplacian to ensure that one has an energy interval around 0 such that for all these energies there are only has elliptic channels.
- If you want to treat an energy interval for the honest Anderson model in a limit to infinite width or if you want to treat all energies in the spectrum of  $H_0$  in this critical limit, then one has to deal with hyperbolic channels.

## Applications in the pure one-dimensional model:

- Kritchevski, Valko and Virag studied the eigenvalue statistics in this critical scaling.
- Rifkind and Virag studied distribution of shape of eigenfunctions in this scaling limit

## Application on strip models:

- Valko and Virag obtained GOE limiting statistics for certain sequences of modified Anderson models on long boxes
- Bachmann and De Roeck discussed relations from Random Matrix Theory to the Anderson model and the DMPK equation
- In both papers the Anderson model on a strip is slightly modified by scaling down the vertical Laplacian to ensure that one has an energy interval around 0 such that for all these energies there are only has elliptic channels.
- If you want to treat an energy interval for the honest Anderson model in a limit to infinite width or if you want to treat all energies in the spectrum of  $H_0$  in this critical limit, then one has to deal with hyperbolic channels.

## Applications in the pure one-dimensional model:

- Kritchevski, Valko and Virag studied the eigenvalue statistics in this critical scaling.
- Rifkind and Virag studied distribution of shape of eigenfunctions in this scaling limit

## Application on strip models:

- Valko and Virag obtained GOE limiting statistics for certain sequences of modified Anderson models on long boxes
- Bachmann and De Roeck discussed relations from Random Matrix Theory to the Anderson model and the DMPK equation
- In both papers the Anderson model on a strip is slightly modified by scaling down the vertical Laplacian to ensure that one has an energy interval around 0 such that for all these energies there are only has elliptic channels.
- **If you want to treat an energy interval for the honest Anderson model in a limit to infinite width or if you want to treat all energies in the spectrum of  $H_0$  in this critical limit, then one has to deal with hyperbolic channels.**

# Heuristics

$X_{\lambda,n} = (\mathcal{T}_0^E)^{-n} \mathcal{T}_{\lambda,[1,n]}^E$  follows an evolution of the form:

$$X_{\lambda,n+1} = \left( \mathbf{1} + \lambda (\mathcal{T}_0^E)^{-n-1} \mathcal{V}_n (\mathcal{T}_0^E)^n \right) X_{\lambda,n}$$

- The  $\mathcal{V}_n$  are i.i.d. random matrices, giving a diffusive term of order  $\lambda^2$  (variance)
- Since  $\mathcal{T}_0^E$  is conjugated to a unitary, i.e. generates a compact group, the conjugations with  $(\mathcal{T}_0^E)^n$  lead to an average over the compact group generated by  $\mathcal{T}_0^E$ .
- In the scaling limit  $n \sim \lambda^{-2} \rightarrow \infty$  the total diffusion after  $n$  steps is of order 1 and one obtains a limiting process as in the central limit theorem.
- If  $\mathcal{T}_0^E$  has eigenvalues of different size, then conjugates lead to exponential growing terms in  $n$ , preventing the existence of a limit process.
- One needs to project on the elliptic channels in a good way.

# Heuristics

$X_{\lambda,n} = (\mathcal{T}_0^E)^{-n} \mathcal{T}_{\lambda,[1,n]}^E$  follows an evolution of the form:

$$X_{\lambda,n+1} = \left( \mathbf{1} + \lambda (\mathcal{T}_0^E)^{-n-1} \mathcal{V}_n (\mathcal{T}_0^E)^n \right) X_{\lambda,n}$$

- The  $\mathcal{V}_n$  are i.i.d. random matrices, giving a diffusive term of order  $\lambda^2$  (variance)
- Since  $\mathcal{T}_0^E$  is conjugated to a unitary, i.e. generates a compact group, the conjugations with  $(\mathcal{T}_0^E)^n$  lead to an average over the compact group generated by  $\mathcal{T}_0^E$ .
- In the scaling limit  $n \sim \lambda^{-2} \rightarrow \infty$  the total diffusion after  $n$  steps is of order 1 and one obtains a limiting process as in the central limit theorem.
- If  $\mathcal{T}_0^E$  has eigenvalues of different size, then conjugates lead to exponential growing terms in  $n$ , preventing the existence of a limit process.
- One needs to project on the elliptic channels in a good way.

# Heuristics

$X_{\lambda,n} = (\mathcal{T}_0^E)^{-n} \mathcal{T}_{\lambda,[1,n]}^E$  follows an evolution of the form:

$$X_{\lambda,n+1} = \left( \mathbf{1} + \lambda (\mathcal{T}_0^E)^{-n-1} \mathcal{V}_n (\mathcal{T}_0^E)^n \right) X_{\lambda,n}$$

- The  $\mathcal{V}_n$  are i.i.d. random matrices, giving a diffusive term of order  $\lambda^2$  (variance)
- Since  $\mathcal{T}_0^E$  is conjugated to a unitary, i.e. generates a compact group, the conjugations with  $(\mathcal{T}_0^E)^n$  lead to an average over the compact group generated by  $\mathcal{T}_0^E$ .
- In the scaling limit  $n \sim \lambda^{-2} \rightarrow \infty$  the total diffusion after  $n$  steps is of order 1 and one obtains a limiting process as in the central limit theorem.
- If  $\mathcal{T}_0^E$  has eigenvalues of different size, then conjugates lead to exponential growing terms in  $n$ , preventing the existence of a limit process.
- One needs to project on the elliptic channels in a good way.



# Heuristics

$X_{\lambda,n} = (\mathcal{T}_0^E)^{-n} \mathcal{T}_{\lambda,[1,n]}^E$  follows an evolution of the form:

$$X_{\lambda,n+1} = \left( \mathbf{1} + \lambda (\mathcal{T}_0^E)^{-n-1} \mathcal{V}_n (\mathcal{T}_0^E)^n \right) X_{\lambda,n}$$

- The  $\mathcal{V}_n$  are i.i.d. random matrices, giving a diffusive term of order  $\lambda^2$  (variance)
- Since  $\mathcal{T}_0^E$  is conjugated to a unitary, i.e. generates a compact group, the conjugations with  $(\mathcal{T}_0^E)^n$  lead to an average over the compact group generated by  $\mathcal{T}_0^E$ .
- In the scaling limit  $n \sim \lambda^{-2} \rightarrow \infty$  the total diffusion after  $n$  steps is of order 1 and one obtains a limiting process as in the central limit theorem.
- If  $\mathcal{T}_0^E$  has eigenvalues of different size, then conjugates lead to exponential growing terms in  $n$ , preventing the existence of a limit process.
- One needs to project on the elliptic channels in a good way.

# Case with hyperbolic channels

- Let us assume we have  $d_e > 0$  elliptic and  $d_h > 0$  hyperbolic channels at energy  $E$ ,  $d_e + d_h = d$ .
- Then with an adequate basis change we find that  $\mathcal{T}_{\lambda,n} = \mathcal{C}\mathcal{T}_{\lambda,n}^E\mathcal{C}^{-1}$  is of the form

$$\mathcal{T}_{\lambda,n} = \begin{pmatrix} \Upsilon & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & U & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \Upsilon^{-1} \end{pmatrix} + \lambda\mathcal{V}(n)$$

where  $U \in U(2d_e)$  is unitary and  $\Upsilon \in GL(d_h)$  satisfies  $\|\Upsilon\| < 1$ ,  $\mathcal{V}(n)$  are i.i.d. random matrices.

- Let

$$\mathcal{X}_{\lambda,n} = \begin{pmatrix} \mathbf{1} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & U^{-n} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{1} \end{pmatrix} \mathcal{T}_{\lambda,n}\mathcal{T}_{\lambda,n-1}\cdots\mathcal{T}_{\lambda,2}\mathcal{T}_{\lambda,1}.$$

We will eliminate the exponential growing part of  $\mathcal{X}_{\lambda,n}$  by taking a Schur complement.

# Case with hyperbolic channels

- Let us assume we have  $d_e > 0$  elliptic and  $d_h > 0$  hyperbolic channels at energy  $E$ ,  $d_e + d_h = d$ .
- Then with an adequate basis change we find that  $\mathcal{T}_{\lambda,n} = \mathcal{C}\mathcal{T}_{\lambda,n}^E\mathcal{C}^{-1}$  is of the form

$$\mathcal{T}_{\lambda,n} = \begin{pmatrix} \Upsilon & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & U & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \Upsilon^{-1} \end{pmatrix} + \lambda\mathcal{V}(n)$$

where  $U \in \mathbf{U}(2d_e)$  is unitary and  $\Upsilon \in \mathbf{GL}(d_h)$  satisfies  $\|\Upsilon\| < 1$ ,  $\mathcal{V}(n)$  are i.i.d. random matrices.

- Let

$$\mathcal{X}_{\lambda,n} = \begin{pmatrix} \mathbf{1} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & U^{-n} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{1} \end{pmatrix} \mathcal{T}_{\lambda,n}\mathcal{T}_{\lambda,n-1}\cdots\mathcal{T}_{\lambda,2}\mathcal{T}_{\lambda,1}.$$

We will eliminate the exponential growing part of  $\mathcal{X}_{\lambda,n}$  by taking a Schur complement.

- Let

$$\mathcal{X}_{\lambda,n} = \begin{pmatrix} \mathbf{1} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & U^{-n} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{1} \end{pmatrix} \mathcal{T}_{\lambda,n} \cdots \mathcal{T}_{\lambda,1} = \begin{pmatrix} A_{\lambda,n} & B_{\lambda,n} \\ C_{\lambda,n} & D_{\lambda,n} \end{pmatrix} \mathcal{X}_0$$

where  $A_{\lambda,n} \in \text{Mat}(d_h + 2d_e, \mathbb{C})$  and  $D_{\lambda,n} \in \text{Mat}(d_h, \mathbb{C})$ .

- Equivalence relation: Let  $\mathcal{X}_1 \sim \mathcal{X}_2$  if  $\mathcal{X}_1 = \mathcal{X}_2 \begin{pmatrix} \mathbf{1}_{2d_e+d_h} & \mathbf{0} \\ C & D \end{pmatrix}$ .
- Since

$$\begin{aligned} & \begin{pmatrix} A_{\lambda,n} & B_{\lambda,n} \\ C_{\lambda,n} & D_{\lambda,n} \end{pmatrix} \begin{pmatrix} \mathbf{1} & \mathbf{0} \\ -D_{\lambda,n}^{-1}C_{\lambda,n} & D_{\lambda,n}^{-1} \end{pmatrix} \\ &= \begin{pmatrix} A_{\lambda,n} - B_{\lambda,n}D_{\lambda,n}^{-1}C_{\lambda,n} & B_{\lambda,n}D_{\lambda,n}^{-1} \\ \mathbf{0} & \mathbf{1}_{d_h} \end{pmatrix} \end{aligned}$$

the equivalence class of  $\mathcal{X}_{\lambda,n}$  is determined by

$$X_{\lambda,n} := A_{\lambda,n} - B_{\lambda,n}D_{\lambda,n}^{-1}C_{\lambda,n}, \quad B_{\lambda,n}D_{\lambda,n}^{-1}$$

- Let

$$\mathcal{X}_{\lambda,n} = \begin{pmatrix} \mathbf{1} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & U^{-n} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{1} \end{pmatrix} \mathcal{T}_{\lambda,n} \cdots \mathcal{T}_{\lambda,1} = \begin{pmatrix} A_{\lambda,n} & B_{\lambda,n} \\ C_{\lambda,n} & D_{\lambda,n} \end{pmatrix} \mathcal{X}_0$$

where  $A_{\lambda,n} \in \text{Mat}(d_h + 2d_e, \mathbb{C})$  and  $D_{\lambda,n} \in \text{Mat}(d_h, \mathbb{C})$ .

- Equivalence relation: Let  $\mathcal{X}_1 \sim \mathcal{X}_2$  if  $\mathcal{X}_1 = \mathcal{X}_2 \begin{pmatrix} \mathbf{1}_{2d_e+d_h} & \mathbf{0} \\ C & D \end{pmatrix}$ .
- Since

$$\begin{aligned} & \begin{pmatrix} A_{\lambda,n} & B_{\lambda,n} \\ C_{\lambda,n} & D_{\lambda,n} \end{pmatrix} \begin{pmatrix} \mathbf{1} & \mathbf{0} \\ -D_{\lambda,n}^{-1}C_{\lambda,n} & D_{\lambda,n}^{-1} \end{pmatrix} \\ &= \begin{pmatrix} A_{\lambda,n} - B_{\lambda,n}D_{\lambda,n}^{-1}C_{\lambda,n} & B_{\lambda,n}D_{\lambda,n}^{-1} \\ \mathbf{0} & \mathbf{1}_{d_h} \end{pmatrix} \end{aligned}$$

the equivalence class of  $\mathcal{X}_{\lambda,n}$  is determined by

$$X_{\lambda,n} := A_{\lambda,n} - B_{\lambda,n}D_{\lambda,n}^{-1}C_{\lambda,n}, \quad B_{\lambda,n}D_{\lambda,n}^{-1}$$

The equivalence class of  $\mathcal{X}_{\lambda,n}$  is determined by the pair

$$X_{\lambda,n} := A_{\lambda,n} - B_{\lambda,n} D_{\lambda,n}^{-1} C_{\lambda,n}, \quad B_{\lambda,n} D_{\lambda,n}^{-1}$$

### Theorem (Virag, S.)

In distribution, for  $n \rightarrow \infty$  and any  $t > 0$ ,

$$B_{\frac{1}{\sqrt{n}}, [tn]} D_{\frac{1}{\sqrt{n}}, [tn]}^{-1} \Longrightarrow \mathbf{0}, \quad X_{\frac{1}{\sqrt{n}}, [tn]} \Longrightarrow X_t = \begin{pmatrix} \mathbf{0} & \mathbf{0} \\ \Lambda_t X_{21} & \Lambda_t X_{22} \end{pmatrix}$$

where  $\Lambda_t$  is a  $2d_e \times 2d_e$  matrix process satisfying a SDE

$$d\Lambda_t = V \Lambda_t dt + dB_t \Lambda_t, \quad \Lambda_0 = \mathbf{1}$$

Doing the same procedure for the transfer matrices  $\mathcal{T}_{\lambda,n}^{E+\lambda^2\varepsilon}$  we obtain the following:

### Theorem (Virag, S.)

Let  $H_{\lambda,n}$  be the restriction of  $H_\lambda$  to  $\ell^2(\{1, \dots, n\}) \otimes \mathbb{C}^d$  with Dirichlet boundary conditions, let  $\mathcal{E}_{\lambda,n}$  be the eigenvalue process of  $H_{\lambda,n} - E$ , then for subsequences  $n_k \rightarrow \infty$

$$n_k \mathcal{E}_{\frac{1}{\sqrt{n_k}}, n_k} \implies \text{zeros}_\varepsilon \det(M_0^* \Lambda_1^\varepsilon M_1)$$

where  $M_0, M_1 \in \mathbb{C}^{2d_e \times d_e}$ ,  $\Lambda_t^\varepsilon$  is a  $2d_e \times 2d_e$  matrix process that for fixed  $\varepsilon$  satisfies a SDE of the form

$$d\Lambda_t^\varepsilon = (V + \varepsilon W)\Lambda_t^\varepsilon dt + dB_t \Lambda_t^\varepsilon, \quad \Lambda_t^\varepsilon = \mathbf{1}$$

THANK YOU!!