## Topological Bifurcations in the Wake Behind an Oscillating Cylinder

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INDEPENDENT
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## von Kármán vortex street

$2 S$ wake



## Oscillating cylinder may lead to exotic wakes

$P+S$ wake


Williamson, in Ponta \& Aref, J. Fluids Struct. 22(2006), 327-344

## Other body shapes produce very exotic wakes



Schnipper et al, J. Fluid Mech. 633(2009), 411-423

## Forced transverse oscillations of the cylinder



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Dimensionless parameters

- $R e=\frac{U D}{\nu}(=100$ here $)$
- $A \leftarrow \frac{A}{D}$
- $f \leftarrow \frac{f}{f_{s t}}=\frac{T_{s t}}{T}$ or

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\lambda \leftarrow \frac{\lambda}{D}=\frac{U T_{s t}}{D}
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where $f_{s t}\left(T_{s t}\right)$ is the frequency (period) of the vortex shedding for $A=0$

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Williamson \& Roshko,
J. Fluids Struct. 2(1988), 355-381


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To elucidate the transitions in the structure of the wake from a $2 S$ to a $\mathrm{P}+\mathrm{S}$ pattern as the forcing amplitude $A$ is varied, at $f=f_{s t}$ and $R e=100$.

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Leontini et al. (2006)

## Locating the vortices, their creation, and disappearance

- $A=1.092:$ note the symmetry $\omega(x, y, t)=-\omega(x,-y, t+1 / 2)$



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Vortices:

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& \partial_{x} \omega\left(x^{*}, y^{*}, t^{*} ; A\right)=0 \\
& \partial_{y} \omega\left(x^{*}, y^{*}, t^{*} ; A\right)=0 \\
& \left|\mathbf{H}^{\omega}\left(x^{*}, y^{*}, t^{*} ; A\right)\right|>0
\end{aligned}
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Saddle points:

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An extremum and a saddle are created or disappear in a cusp bifurcation at a degenerate critical point of vorticity where the Hessian is singular.


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This wake is classified as $(2 P)^{3}(2 S)^{\infty}$

- Extended classification concatenated by symbols $\Sigma^{n}$ with $\Sigma=2 \mathrm{~S}, \mathrm{P}+\mathrm{S}, 2 \mathrm{P}, \ldots$.
- Integer $n$ designates the number of periods +1 the pattern exists.


## Bistability

- $A=1.092$

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## Bistability

- $\omega(x, y, t)=-\omega(x,-y, t+1 / 2)$

- $\omega(x, y, t) \neq-\omega(x,-y, t+1 / 2)$



## A dynamical subcritical bifurcation

Matharu, Hazel \& Heil, J. Fluid Mech. 918(2021), A42:

- A subcritical symmetry-breaking pitchfork bifurcation from the symmetric branch occurs at $A_{2}=1.093$
- The bifurcating branch gains stability at a fold bifurcation at $A_{1}=1.078$

$$
\varepsilon=\|\omega(x, y, t)+\omega(x,-y, t+1 / 2)\|
$$



## Topological bifurcation diagrams




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$$
A<A_{0}:(2 S)^{\infty}=2 S
$$

## Topological bifurcation diagrams




$$
A_{0}<A<A_{2}:(2 S)^{n}(2 \mathrm{P})^{m}(2 \mathrm{~S})^{\infty}
$$

## Topological bifurcation diagrams



$A_{2}<A<A_{3}:(\mathrm{P}+\mathrm{S})^{n}(2 \mathrm{P})^{m}(\mathrm{P}+\mathrm{S})^{\infty}$

## Topological bifurcation diagrams



$A_{3}<A:(\mathrm{P}+\mathrm{S})^{\infty}=\mathrm{P}+\mathrm{S}$

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$A_{3}<A:(\mathrm{P}+\mathrm{S})^{\infty}=\mathrm{P}+\mathrm{S}$



## Following the unstable branch



## Varying the forcing period




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- A sequence of intermediate patterns of the form $(2 S)^{n}(2 P)^{m}(2 S)^{\infty}$ and $(\mathrm{P}+\mathrm{S})^{n}(2 \mathrm{P})^{m}(\mathrm{P}+\mathrm{S})^{\infty}$ occur in the transition process.


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- Jumps in the wake structure occur at the subcritical dynamical pitchfork bifurcation points $A_{1}$ and $A_{2}$ of the periodic flow, $A_{0}<A_{1}<A_{2}<A_{3}$. The dynamical pitchfork bifurcation is required to obtain asymmetric patterns $P+S$


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- The presence of the (overlooked) local patterns (2P) ${ }^{m}$ may be related to the mystery of the missing 2P wake.
Details in AR Nielsen, PS Matharu \& M Brøns: Topological bifurcations in the transition from two single vortices to a pair and a single vortex in the periodic wake behind an oscillating cylinder, J. Fluid Mech. 940(2022), A22.

